

On the Nonlocal Cahn-hilliard Inpainting Model with Singular Nonlinearity

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Abstract: In this paper, we consider the nonlocal Cahn-Hilliard equation with a fidelity term and singular nonlinearities, which has numerous applications in biology and image inpainting. First, we prove the existence of a unique solution. We approximate the singular nonlinearities by regular nonlinearities, from which we derive important priori-estimates for the solution of the approximating problem. We then establish the existence of a solution to the approximating problem using the Faedo-Galerkin method. Additionally, we demonstrate that the solution of the Cahn-Hilliard equation in question is of higher regular by applying Sobolev embedding theorems. Finally, we discuss the important separation property of the solution that allows us to obtain a global (in time) unique solution of the problem.

Key- Words: Nonlocal Cahn-Hilliard, Well-posedness, Singular potential, Image Inpainting, Separation property, Galerkin algorithm

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1 Introduction

The Cahn-Hilliard equation, proposed in 1958, describes phase separation between two materials. This model is the natural choice from a mathematical point of view. However, from a physical (microscopic) perspective, when the domain is too large, the ε -neighborhood in the classical model becomes insufficient. To address this, Giacomini and Lebowitz proposed a nonlocal model in 1992 to describe the same phenomenon from a microscopic standpoint. Subsequently, Bates and Han proposed a modification of the Giacomini-Lebowitz model. We refer the readers for more details [1], [2], [3], [4], [5], [6], [7], [8], [9], [10], [11], [12], [13], [14], [15], [16], [17], [19], [20], [21].

Giacomini and Lebowitz pointed out that this equation cannot be rigorously derived from microscopic principles. Using statistical

mechanics and the hydrodynamic limit, they formulated a continuum model—a nonlocal Cahn-Hilliard equation-governed by the following free energy functional:

$$\mathcal{E}_{nCH1} = -\frac{1}{2} \int_{\Omega} \int_{\Omega} M(x-y) \varphi(x) \varphi(y) dx dy + \int_{\Omega} S(\varphi(x)) dx.$$

This equation is known as Helmholtz free energy, and $M : \mathbb{R}^l \rightarrow \mathbb{R}^l$ represents a smooth convolution kernel, fulfilling $M(x) = M(-x)$. Additionally, S is a convex potential, typically approximated by a convex polynomial and defined as:

$$S(m) = (m+1) \ln(m+1) + (1-m) \ln(1-m), \quad m \in]-1, 1[. \quad (1)$$

Then, the nonlocal Cahn-Hilliard equation given

by

$$\frac{\partial \varphi}{\partial t} = \Delta v, \quad v := \frac{\partial \mathcal{E}_{nCH1}}{\partial \varphi} = f(\varphi) - M *_{\Omega} \varphi, \quad (2)$$

where

$$(M *_{\Omega} \varphi)(x, t) = \int_{\Omega} M(x-y) \varphi(y, t) dy, \quad x \in \Omega, t \geq 0.$$

The nonlocal Cahn-Hilliard energy can also be expressed in the form

$$\begin{aligned} \mathcal{E}_{nCH2} = & \frac{1}{4} \int_{\Omega} \int_{\Omega} M(x-y) (\varphi(x) - \varphi(y))^2 dx dy \\ & + \int_{\Omega} F(\varphi(x)) dx, \end{aligned}$$

was proposed by the authors Bates and Han in [3,4]. Here, F denotes the double-well potential, as in the classical Cahn-Hilliard model, and is a non-convex function. It is typically approximated by a smooth potential of the following form:

$$F(m) = \frac{m^4}{4} + \frac{\lambda_1 - \lambda_2}{2} m^2, \quad \forall m \in \mathbb{R}, \quad (3)$$

where $\lambda_j, j = 1, 2$ are positive constants. Otherwise, if $\lambda_1 - \lambda_2 > 0$, then F can be seen as a smooth approximation of S .

1.1 Notations

Let $\Omega \subset \mathbb{R}^l$ be bounded, and Y a real Banach space with dual Y^* . We also introduce the following two spaces:

$$H^1(0, T, Y, Y^*) = \{\rho \in L^2(0, T, Y); \rho_t \in L^2(0, T, Y^*)\}$$

and

$$H_0^1(\Omega) = \{\rho \in H^1(\Omega); \bar{\rho} = 0\},$$

where $\bar{\rho} = \langle \rho, \frac{1}{|\Omega|} \rangle \quad \forall \rho \in (H^1(\Omega))^*$.

The inner product $\langle \cdot, \cdot \rangle$ represents the pairing of $H^1(\Omega)$ and $(H^1(\Omega))^*$ and the variables ρ' , $\partial_t \rho$, $\frac{\partial \rho}{\partial t}$, and ρ_t represent the partial derivative of ρ with respect to time t .

Lastly, for any $\varrho \in H^1([0, T] \times \Omega)$, we denote its partial derivative with respect to time by $\partial_t \varrho$ or $\frac{\partial \varrho}{\partial t}$ and its partial derivative with respect to the x_j -variable by $\partial_j \varrho$.

Suppose $\lambda : \mathbb{R}^N \rightarrow \mathbb{R}$ and $\gamma : \mathbb{R}^N \rightarrow \mathbb{R}$ are measurable functions. Then, the convolution of λ and γ defined as :

$$\lambda * \gamma(x) = \int_{\Omega} \lambda(x-y) \gamma(y) dy \text{ for } x \in \mathbb{R}^N$$

The values of C, R and E may vary, even within the same line.

1.2 Assumptions

Consider the following assumptions:

(A) The convolution kernel $M : \mathbb{R}^l \rightarrow \mathbb{R}$ satisfies the following conditions:

$$M(x) = M(-x), \text{ a.a. } x \in \mathbb{R}^l, \quad (A1)$$

$$\sup_{x \in \Omega} \int_{\Omega} |M(x-y)| < +\infty, \quad (A2)$$

$$\begin{aligned} \forall p \in [1, +\infty], \exists r_p > 0 \text{ such that} \\ \|M * \vartheta\|_{W^{1,p}(\Omega)} \leq r_p \|\vartheta\|_{L^p(\Omega)}, \end{aligned} \quad (A3)$$

$$\exists R > 0 \text{ such that } \|M * \vartheta\|_{W^{2,2}(\Omega)} \leq R \|\vartheta\|_{W^{1,2}(\Omega)}. \quad (A4)$$

(B) The initial condition ϕ_0 is assumed to satisfy

$$\phi_0 \text{ is measurable}, \quad (B1)$$

$$0 \leq \phi_0(x) \leq 1 \text{ a.e. } x \in \Omega, \quad (B2)$$

$$0 < \bar{\phi}_0 < 1. \quad (B3)$$

(C) The damaged image $h \in L^2(\Omega)$ takes values between 0 and 1.

(D) The constant λ_0 is positive and assumed to be large.

(E) The indicator function $\chi_{\Omega \setminus D}(x)$ is defined by:

$$\chi_{\Omega \setminus D}(x) = \begin{cases} 0 & , x \in D \\ 1 & , x \in \Omega \setminus D. \end{cases}$$

2 Problem Formulation

We study the initial-boundary value problem given by:

$$\partial_t \varphi - \nabla \cdot (\mu \nabla v) = \lambda_0 \chi_{\Omega \setminus D}(x) (h - \varphi), \text{ in } \Omega \times (0, T), \quad (4)$$

$$v = f'(\varphi) + \theta, \text{ in } \Omega \times (0, T), \quad (5)$$

$$\theta(x, t) = \int_{\Omega} M(|x-y|) (1-2\varphi(y, t)) dy \text{ in } \Omega \times (0, T), \quad (6)$$

$$n \cdot \mu \nabla v = 0, \text{ on } \partial \Omega \times (0, T), \quad (7)$$

$$\varphi(0) = \varphi_0 \text{ in } \Omega, \quad (8)$$

where $\lambda_0 \chi_{\Omega \setminus D}(x) (h - \varphi)$ is a fidelity term used in image inpainting.

The function $h(x)$ represents the damaged image, D denotes the inpainting domain, and λ_0 is a constant. Equations (1.7)–(1.11) correspond to the nonlocal Cahn-Hilliard equation with a fidelity term and singular nonlinearities. Moreover, n is the outward unit normal to $\partial \Omega$, and M is the convolution kernel. The functions f and μ are defined as:

$$f(\varphi) = \varphi \log(\varphi) + (1 - \varphi) \log(1 - \varphi),$$

$$\mu = \frac{1}{f''(\varphi)} = \varphi(1 - \varphi),$$

where μ is the degenerate mobility.

The existence, uniqueness, regularity, and separation properties of the nonlocal Cahn-Hilliard equation with a fidelity term and singular nonlinearities have not been established in previous works. However, the authors in [12] proved the well-posedness of the convective nonlocal Cahn-Hilliard equation with a reaction term. Furthermore, in [?], the existence and uniqueness of the nonlocal Cahn-Hilliard equation with a reaction term are established.

3 Problem Solution

We consider the nonlocal Cahn-Hilliard equation with a fidelity term, singular (logarithmic) nonlinearities, and degenerate mobility (i.e., non-constant) in the form:

$$\partial_t \varphi - \nabla \cdot (\mu(\nabla v)) = \lambda_0 \chi_{\Omega \setminus D}(x)(h - \varphi), \text{ in } \psi, \quad (9)$$

$$v = f'(\varphi) + \theta, \text{ in } \psi, \quad (10)$$

$$\theta(x, t) = \int_{\Omega} M(|x - y|)(1 - 2\varphi(y, t))dy, \text{ in } \psi, \quad (11)$$

$$n \cdot \mu \nabla v = 0, \text{ on } \Psi, \quad (12)$$

$$\varphi(0) = \varphi_0, \text{ in } \Omega, \quad (13)$$

where $\psi = \Omega \times (0, T)$, $\Omega \subseteq \mathbb{R}^l$, and $\Psi = \partial\Omega \times (0, T)$. Furthermore, the functions f , and μ are defined as:

$$f(\varphi) = \varphi \ln \varphi + (1 - \varphi) \ln(1 - \varphi), \quad (14)$$

$$\mu = \frac{1}{f''(\varphi)} = \varphi(1 - \varphi). \quad (15)$$

Definition 3.1 Suppose that conditions (A1)-(A4) and (B1)-(B3) hold. Let $T \in (0, \infty)$. If

$$\varphi \in H^1(0, T, H^1(\Omega), (H^1(\Omega))^*), \quad (16)$$

$$0 \leq \varphi \leq 1 \text{ a.e in } \psi, \quad (17)$$

$$\theta = M * (1 - 2\varphi) \text{ a.e. in } \psi,$$

$$\theta \in C([0, T], W^{1,\infty}(\Omega)),$$

$$\varphi(0) = \varphi_0 \text{ in } L^2(\Omega),$$

then, φ is a weak solution of (9)-(13) on $[0, T]$,

Theorem 3.2 Under the above assumptions, there exists a unique weak solution of (??)-(??) (in the sense of the preceding definition) satisfying

$$\varphi \in H^1(0, T, H^1(\Omega), (H^1(\Omega))^*) (\hookrightarrow C([0, T], L^2(\Omega)))$$

3.0.1 Uniqueness

We proceed to establish the uniqueness of the solution.

Proof: Let φ_i , $i \in \{1, 2\}$, be two solutions of the system (9)-(13) with corresponding initial data $\varphi_{0,i}$, $i \in \{1, 2\}$. Our goal is to prove that the solution is unique, i.e., $\varphi_1 = \varphi_2$.

The variational formulation of equations (9)-(10) reads as follows:

We multiply (9)-(10) by a test function $\phi \in H^1(\psi)$ respectively and integrate over Ω :

$$((\partial_t \varphi, \phi)) = -((\nabla \cdot \mu(\nabla v), \nabla \phi)) + ((\chi_{\Omega \setminus D}(x)(h - \varphi), \phi)), \quad (18)$$

$$((v, \nabla \phi)) = ((f'(\varphi), \nabla \phi)) + ((\theta, \nabla \phi)). \quad (19)$$

Using (15), (18), and (19), we get for $i \in \{1, 2\}$:

$$((\partial_t \varphi_i, \phi)) = -((\nabla \varphi_i + \mu_i \nabla \theta_i, \nabla \phi)) + ((\chi_{\Omega \setminus D}(x)(h - \varphi_i), \phi)) \quad \forall \phi \in H^1(\psi), \text{ a.e in } (0, T), \quad (20)$$

$$\mu_i = \mu(\varphi_i) = \varphi_i(1 - \varphi_i), \text{ for } i \in \{1, 2\}, \quad (21)$$

$$\theta_i = M * (1 - 2\varphi_i), \text{ for } i \in \{1, 2\}. \quad (22)$$

Let ϕ be the difference between the two solutions, i.e. set $\phi = \varphi = \varphi_1 - \varphi_2$.

For each $i \in \{1, 2\}$, we have:

$$((\partial_t \varphi_1, \phi)) = -((\nabla \varphi_1 + \mu_1 \nabla \theta_1, \nabla \phi)) + ((\chi_{\Omega \setminus D}(x)(h - \varphi_1), \phi)), \quad (23)$$

$$((\partial_t \varphi_2, \phi)) = -((\nabla \varphi_2 + \mu_2 \nabla \theta_2, \nabla \phi)) + ((\chi_{\Omega \setminus D}(x)(h - \varphi_2), \phi)). \quad (24)$$

By subtracting now equations (23) and (24), we obtain:

$$((\partial_t \varphi, \phi)) = -((\nabla \varphi, \nabla \phi)) - ((\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2, \nabla \phi)) + ((\chi_{\Omega \setminus D}(x)(h - \varphi_1) - \chi_{\Omega \setminus D}(x)(h - \varphi_2), \phi)). \quad (25)$$

Furthermore, equation (25) is equivalent to:

$$\frac{1}{2} \frac{d}{dt} |\varphi|^2 = - \int_{\Omega} |\nabla \varphi|^2 - \int_{\Omega} (\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2) \nabla \varphi + \int_{\Omega \setminus D} (\chi_{\Omega \setminus D}(x)(h - \varphi_1) - \chi_{\Omega \setminus D}(x)(h - \varphi_2)) \varphi. \quad (26)$$

Then, integrating equation (26) over $(0, t)$, $t \in (0, T]$ and using the fact that

$$\int_0^t \frac{d}{dt} \varphi = \varphi(t) - \varphi_0,$$

we get

$$\begin{aligned} & \frac{1}{2} \|\varphi(t)\|_{L^2(\Omega)}^2 - \frac{1}{2} \|\varphi_{01} - \varphi_{02}\|_{L^2(\Omega)}^2 \\ &= - \int_0^t \int_{\Omega} |\nabla \varphi|^2 - \int_0^t \int_{\Omega} (\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2) \nabla \varphi \\ &+ \int_0^t \int_{\Omega \setminus D} (\chi_{\Omega \setminus D}(x)(h - \varphi_1) - \chi_{\Omega \setminus D}(x)(h - \varphi_2)) \varphi \\ &\leq - \int_0^t \int_{\Omega} |\nabla \varphi|^2 + \int_0^t \int_{\Omega} |\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2| |\nabla \varphi| \\ &+ \int_0^t \int_{\Omega \setminus D} |\chi_{\Omega \setminus D}(x)(h - \varphi_1) - \chi_{\Omega \setminus D}(x)(h - \varphi_2)| |\varphi|. \end{aligned} \quad (27)$$

Finally, we obtain the following estimates by applying Young's inequality:

$$\begin{aligned} & \int_{\Omega} |\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2| |\nabla \varphi| \leq \frac{1}{2} \int_{\Omega} |\nabla \varphi|^2 \\ &+ \frac{1}{2} \int_{\Omega} |\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2|^2, \end{aligned} \quad (*)$$

$$\begin{aligned} & \int_{\Omega} |\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2|^2 \\ &= \int_{\Omega} |\mu_1 \nabla \theta_1 - \mu_2 \nabla \theta_2 - \mu_1 \nabla \theta_2 + \mu_1 \nabla \theta_2|^2 \\ &\leq \int_{\Omega} |\mu_1|^2 |\nabla \theta_1 - \nabla \theta_2|^2 + \int_{\Omega} |\nabla \theta_2|^2 |\mu_1 - \mu_2|^2 \\ &\leq C \|\nabla \theta_1 - \nabla \theta_2\|_{L^2(\Omega)}^2 + \int_{\Omega} |\varphi(1 - \varphi_1 - \varphi_2)|^2 |\nabla \theta_2|^2 \\ &\leq C \|\nabla \theta_1 - \nabla \theta_2\|_{L^2(\Omega)}^2 + C \int_{\Omega} |\varphi|^2 |\nabla \theta_2|^2 \\ &\leq C \|\nabla \theta_1 - \nabla \theta_2\|_{L^2(\Omega)}^2 + C \|\nabla \theta_2\|_{L^\infty(\Omega)}^2 \|\varphi\|_{L^2(\Omega)}^2 \\ &\leq C 4r_2^2 \|\varphi\|_{L^2(\Omega)}^2 + C 4r_\infty^2 \|\varphi\|_{L^2(\Omega)}^2 \\ &\leq C \|\varphi\|_{L^2(\Omega)}^2, \end{aligned} \quad (**)$$

$$\begin{aligned} & \int_{\Omega \setminus D} |\chi_{\Omega \setminus D}(h - \varphi_1) - \chi_{\Omega \setminus D}(h - \varphi_2)| \\ &\leq \int_{\Omega} |h - \varphi_1 - (h - \varphi_2)| |\varphi| \\ &\leq \int_{\Omega} |\varphi_1 - \varphi_2| |\varphi| \\ &\leq \int_{\Omega} |\varphi| |\varphi| \\ &\leq \|\varphi\|_{L^2(\Omega)}^2. \end{aligned} \quad (***)$$

Next, by substituting (*), (**) and (***) into

(3.19), we obtain:

$$\begin{aligned} & \frac{1}{2} \|\varphi(t)\|_{L^2(\Omega)}^2 - \frac{1}{2} \|\varphi_{01} - \varphi_{02}\|_{L^2(\Omega)}^2 \\ &= - \int_0^t \int_{\Omega} |\nabla \varphi|^2 + \frac{1}{2} \int_0^t \int_{\Omega} |\nabla \varphi|^2 \\ &\quad + \frac{1}{2} C \int_0^t \|\varphi\|_{L^2(\Omega)}^2 + \int_0^t \|\varphi\|_{L^2(\Omega)}^2 \\ &\leq - \frac{1}{2} \int_0^t \|\nabla \varphi\|_{L^2(\Omega)}^2 + C \int_0^t \|\varphi\|_{L^2(\Omega)}^2 + \int_0^t \|\varphi\|_{L^2(\Omega)}^2 \\ &\leq C \int_0^t \|\varphi\|_{L^2(\Omega)}^2. \end{aligned} \quad (28)$$

Then, we obtain:

$$\|\varphi(t)\|_{L^2(\Omega)}^2 \leq \|\varphi_{01} - \varphi_{02}\|_{L^2(\Omega)}^2 + C \int_0^t \|\varphi\|_{L^2(\Omega)}^2.$$

Applying Gronwall's Lemma, we reach:

$$\|\varphi(t)\|_{L^2(\Omega)}^2 \leq e^{Ct} \|\varphi_{01} - \varphi_{02}\|_{L^2(\Omega)}^2.$$

Furthermore, since $\varphi_{01} = \varphi_{02}$,

$$\|\varphi(t)\|_{L^2(\Omega)}^2 = 0 \iff \varphi = 0 \iff \varphi_1 = \varphi_2.$$

Hence, the problem (3.1)-(3.5) have a unique solution. $\square$

3.0.2 Existence

We consider the approximate problem P_ε defined as follows ($\forall \in H^1(\Omega)$, a.e in $(0, T)$):

$$((\partial_t \varphi, \phi)) + ((\mu_\varepsilon \nabla v, \nabla \phi)) = ((\chi_{\Omega \setminus D}(h - \varphi), \phi)) \quad (29)$$

$$v = f'_\varepsilon(\varphi) + \theta \text{ a.e in } \psi \quad (30)$$

$$\theta = M * (1 - 2\varphi) \text{ a.e in } \psi \quad (31)$$

$$n \cdot \mu_\varepsilon \nabla v = 0, \text{ on } \Psi, \quad (32)$$

$$\varphi(0) = \varphi_0, \text{ a.e. } x \in \psi, \quad (33)$$

where f'_ε and μ_ε are defined as follows:

$$\mu_\varepsilon = \max\{\mu + \varepsilon, \varepsilon\} \quad (34)$$

The function f'_ε is a solution to the following Cauchy problem:

$$f''_\varepsilon = (1 + 2b_\varepsilon) \frac{1}{\mu_\varepsilon}, \text{ where } b_\varepsilon = \frac{(1 + 4\varepsilon)^{\frac{1}{2}} - 1}{2}, \quad (35)$$

$$f'_\varepsilon\left(\frac{1}{2}\right) = f'\left(\frac{1}{2}\right), f_\varepsilon\left(\frac{1}{2}\right) = f\left(\frac{1}{2}\right). \quad (36)$$

On the other hand, μ_ε can be written as:

$$\mu_\varepsilon(m) = \begin{cases} \varepsilon & \text{for } m < 0, \\ (m + b_\varepsilon)(1 + b_\varepsilon - m) & \text{for } m \in [0, 1], \\ \varepsilon & \text{for } m > 1. \end{cases} \quad (37)$$

From this, we deduce that μ_ε is continuous and

$$\varepsilon \leq \mu_\varepsilon \leq \mu_\varepsilon \left(\frac{1}{2} \right) = \frac{1+4\varepsilon}{4}. \quad (38)$$

Moreover, f''_ε has the following form:

$$f''_\varepsilon(m) = \begin{cases} \frac{1+2b_\varepsilon}{\varepsilon} & \text{for } m < 0, \\ \frac{1+2b_\varepsilon}{(m+b_\varepsilon)(1+b_\varepsilon-m)} & \text{for } m \in [0, 1], \\ \frac{1+2b_\varepsilon}{\varepsilon} & \text{for } m > 1, \end{cases} \quad (39)$$

and

$$0 \leq f''_\varepsilon \leq \frac{1+2w_\varepsilon}{\varepsilon}. \quad (40)$$

Remark 3.3 The function f''_ε possesses the following symmetry property:

$$f''_\varepsilon \left(\frac{1}{2} + m \right) = f''_\varepsilon \left(\frac{1}{2} - m \right) \quad \forall m \in \mathbb{R}.$$

Using the fact that $f'_\varepsilon \left(\frac{1}{2} \right) = f' \left(\frac{1}{2} \right) = 0$, it follows that $f'_\varepsilon(m) < 0$ for $m < \frac{1}{2}$ and $f'_\varepsilon(m) > 0$ for $m > \frac{1}{2}$. Hence,

$$\begin{cases} f'_\varepsilon(m) < 0 & \text{for } m < 0, \\ f'_\varepsilon(m) = \ln \left(\frac{b_\varepsilon+m}{1+b_\varepsilon-m} \right) & \text{for } m \in [0, 1], \\ f'_\varepsilon(m) > 0 & \text{for } m > 1, \end{cases} \quad (41)$$

and

$$f'_\varepsilon \left(\frac{1}{2} + m \right) = -f'_\varepsilon \left(\frac{1}{2} - m \right) \quad \forall m \in \mathbb{R}. \quad (42)$$

Additionally, knowing that $f''_\varepsilon \leq \frac{1+2b_\varepsilon}{\varepsilon}$ and $f'_\varepsilon \left(\frac{1}{2} \right) = 0$, we obtain:

$$|f'_\varepsilon(m)| \leq \frac{1+2b_\varepsilon}{\varepsilon} \left| m - \frac{1}{2} \right| \quad \forall s \in \mathbb{R}. \quad (43)$$

we thus infer that

$$f_\varepsilon \left(\frac{1}{2} + m \right) = f_\varepsilon \left(\frac{1}{2} - m \right) \quad \forall m \in \mathbb{R}. \quad (44)$$

Likewise, for $m > 1$, we have that

$$f'_\varepsilon(m) = \frac{1+2b_\varepsilon}{\varepsilon}(m-1) + f'_\varepsilon(1).$$

However, $f'_\varepsilon(m) \geq 0$ for $m \geq 0$, so that

$$f'_\varepsilon(m) \geq \frac{1+2b_\varepsilon}{\varepsilon}m - \frac{1+2b_\varepsilon}{\varepsilon} \quad \forall m > \frac{1}{2}. \quad (45)$$

Thanks Young's inequality, we get

$$\begin{aligned} & f_\varepsilon(m) - f_\varepsilon \left(\frac{1}{2} \right) \\ & \geq \frac{1+2b_\varepsilon}{2\varepsilon}m^2 - \frac{1+2b_\varepsilon}{\varepsilon}m - \frac{1+2b_\varepsilon}{2\varepsilon}\frac{1}{4} + \frac{1+2b_\varepsilon}{\varepsilon}\frac{1}{2} \\ & \geq \frac{1+2b_\varepsilon}{2\varepsilon}m^2 - \delta m^2 - \frac{1+2b_\varepsilon}{2\varepsilon}\frac{1}{4\delta} - \frac{1+2b_\varepsilon}{2\varepsilon}\frac{1}{4} + \frac{1+2b_\varepsilon}{\varepsilon}\frac{1}{2}, \\ & \forall \delta > 0. \text{ Since } \frac{1+2b_\varepsilon}{2\varepsilon} > \frac{1+b_\varepsilon}{2\varepsilon} \text{ and by choosing } \delta \text{ to be sufficiently small, we get:} \end{aligned}$$

$$f_\varepsilon(m) \geq \frac{1+b_\varepsilon}{2\varepsilon} \left(m - \frac{1}{2} \right)^2 - q'_\varepsilon \quad \forall m \in \mathbb{R}, \quad (46)$$

such that q'_ε is a positive constant that depends on ε .

Besides, we have that

$$\begin{aligned} f_\varepsilon(m) & \geq \frac{1+b_\varepsilon}{2\varepsilon} \left(m^2 - m - \frac{1}{4} \right) - q'_\varepsilon \\ & \geq \frac{1+b_\varepsilon}{2\varepsilon} \left((1-\delta)m^2 - \frac{1}{4} - \frac{1}{8\delta} \right) - q'_\varepsilon, \end{aligned}$$

for all $\delta > 0$. By choosing δ sufficiently small and noting that $\frac{1+b_\varepsilon}{2\varepsilon} > \frac{1}{2\varepsilon}$, it follows that the below inequality holds:

$$f_\varepsilon(m) \geq \frac{1}{2\varepsilon}m^2 - q_\varepsilon \quad \forall m \in \mathbb{R}, \quad (47)$$

where $q_\varepsilon > 0$ depends on ε .

Lemma 3.4 Let assumptions (A) be satisfied for $\varepsilon < \frac{1}{2r_2}$. Then, there exists a function φ such that

$\varphi \in H^1(0, T, H^1(\Omega), (H^1(\Omega))^*) \cap L^\infty(0, T, L^2(\Omega))$ is solution of (29)-(33) with

$$\left\| \mu_\varepsilon^{\frac{1}{2}}(\varphi) |\nabla v| \right\|_{L^2(0, T, L^2(\Omega))} \leq E,$$

where E is a constant that depends of ε .

Proof: The construction of the proof is formed by using Faedo-Galerkin's method.

Let $Y = \{\varphi \in H^1(\Omega); \int_\Omega \varphi dx = 0\}$ be a closed subspace of H^1 .

Then, there exist eigenvalues $\lambda_1, \lambda_2, \dots, \lambda_N$ and $\gamma_1, \gamma_2, \dots, \gamma_N$ with corresponding eigenfunctions $y_1, \dots, y_N$, which form a Galerkin basis for the operator $-\Delta + Id : Y \longrightarrow Y^*$, where $Y \subseteq H^1(\Omega)$.

Consider the space $Y_N = span(\{y_j\}_{j=1}^N)$.

Define the orthogonal projection

$$P_N : H = L^2(\Omega) \longrightarrow Y_N$$

with functions:

$$\varphi_N = \sum_{j=1}^N \lambda_j(t) y_j \quad \text{and} \quad v_N = \sum_{j=1}^N \gamma_j(t) y_j,$$

which solves the following approximate problem:

$$\begin{aligned} & ((\varphi'_N, \phi)) + ((\mu_\varepsilon(\varphi_N) \nabla v_N, \nabla \phi)) \quad (48) \\ & = ((\chi_{\Omega \setminus D}(x)(h - \varphi_N), \phi)) \quad \forall \phi \in Y_N, \\ & v_N = P_N(M * (1 - 2\varphi_N) + f'_\varepsilon(\varphi_N)), \\ & \chi_{\Omega \setminus D}(x)(h - \varphi_N) = P_N(\chi_{\Omega \setminus D}(x)(h - \varphi_N)), \\ & \varphi_N(0) = \varphi_{0N}. \quad (49) \end{aligned}$$

From (34), (35), and (36), we deduce that the function $(\mu_\varepsilon(\varphi_N) \nabla v_N, \nabla \phi)$, for all $\phi \in Y_N$, is locally Lipschitz with respect to the variables λ_j , uniformly in t . This implies that there exists $T_N \in \mathbb{R}_+$ such that problem (48) admits a unique solution $\lambda_1, \dots, \lambda_N, \gamma_1, \dots, \gamma_N \in C^1([0, T_N]; \mathbb{R})$.

Next, we derive a priori estimates to confirm that φ_N is uniformly bounded in N as follows: multiply equation (48) by the test function $\phi = v_N$. Then, obtain:

$$\begin{aligned} & ((\varphi'_N, v_N)) + ((\mu_\varepsilon(\varphi_N) \nabla v_N, \nabla v_N)) \\ & = ((\chi_{\Omega \setminus D}(x)(h - \varphi_N), v_N)). \end{aligned}$$

We now compute the function $((\varphi'_N, v_N))$ as follows:

$$\begin{aligned} & ((\varphi'_N, v_N)) = ((\varphi'_N, f'_\varepsilon(\varphi_N)) + ((\varphi'_N, M * (1 - 2\varphi_N))) \\ & = \int_{\Omega} f'_\varepsilon(\varphi'_N) + \int_{\Omega} M * (1 - 2\varphi_N) \varphi'_N \\ & = \int_{\Omega} \frac{d}{dt} f_\varepsilon(\varphi_N) + \int_{\Omega} \frac{d}{dt} (M * (1 - \varphi_N)) \varphi_N \\ & = \frac{d}{dt} \left(\int_{\Omega} f_\varepsilon(\varphi_N) + \int_{\Omega} \int_{\Omega} M(x - y) \varphi_N(x) (1 - \varphi_N(y)) \right). \quad (50) \end{aligned}$$

Therefore

$$\begin{aligned} & \left(\int_{\Omega} f_\varepsilon(\varphi_N) + \int_{\Omega} \int_{\Omega} M(x - y) \varphi_N(x) (1 - \varphi_N(y)) \right) (t) \\ & \quad + \int_0^t \int_{\Omega} \mu_\varepsilon(\varphi_N) |v_N|^2 \\ & = \int_0^t (\chi_{\Omega \setminus D}(x)(h - \varphi_N), v_N) \\ & \quad + \left(\int_{\Omega} f_\varepsilon(\varphi_N) + \int_{\Omega} \int_{\Omega} M(x - y) \varphi_N(x) (1 - \varphi_N(y)) \right) (0). \end{aligned}$$

Besides, we have $|f'_\varepsilon(m)| \leq C|m| + C$. Thus, by using this estimate and applying assumption (A) along with Young's inequality, we obtain:

$$\begin{aligned} & \int_0^t (\chi_{\Omega \setminus D}(x)(h - \varphi_N), v_N) \\ & = \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(h - \varphi_N) f'_\varepsilon(\varphi_N) \\ & \quad + \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(h - \varphi_N), M * (1 - 2\varphi_N) \\ & \leq \int_{\Omega} (h - \varphi_N)(C|\varphi_N| + C) + \int_{\Omega} [M * (1 - 2\varphi_N)](h - \varphi_N) \\ & \leq C\|\varphi_N\|_H^2 + C\|\varphi_N\|_H + C + r_2\|1 - 2\varphi_N\|_H \|h - \varphi_N\|_H \\ & \leq C\|\varphi_N\|_H^2 + C\|\varphi_N\|_H + C\|\varphi_N\|_H \|h\|_H + C \\ & \leq C\|\varphi_N\|_H^2 + \frac{C}{2}\|\varphi_N\|_H^2 + \frac{C}{2} + \frac{C}{2}\|\varphi_N\|_H^2 + \frac{C}{2}\|h\|_H^2 + C \\ & \leq C\|\varphi_N\|_H^2 + E. \end{aligned}$$

Using (47), assumption (A), and Young's inequality, we get:

$$\begin{aligned} & \int_{\Omega} f_\varepsilon(\varphi_N) + \int_{\Omega} \int_{\Omega} M(x - y) \varphi_N(x) (1 - \varphi_N(y)) \\ & \geq \frac{1}{2\varepsilon} \int_{\Omega} \varphi_N^2 - q_\varepsilon + \int_{\Omega} [M * (1 - \varphi_N)] \varphi_N \\ & \geq \frac{1}{2\varepsilon} \|\varphi_N\|_H^2 - q_\varepsilon - r_2 \|\varphi_N\|_H \|1 - \varphi_N\|_H \\ & \geq \frac{1}{2\varepsilon} \|\varphi_N\|_H^2 - q_\varepsilon - r_2 \|\varphi_N\|_H^2 - r_2 |\Omega| \|\varphi_N\|_H \\ & \geq \left(\frac{1}{2\varepsilon} - r_2 - \delta \right) \|\varphi_N\|_H^2 - q_\varepsilon - \frac{r_2 |\Omega|}{4\delta} \\ & \geq \left(\frac{1}{2\varepsilon} - r_2 - \delta \right) \|\varphi_N\|_H^2 - q_{\delta, \varepsilon} \quad \forall \delta > 0, \quad (51) \end{aligned}$$

where $q_{\delta, \varepsilon}$ denotes a constant that depends on both δ and ε .

Since $\frac{1}{2\varepsilon} - r_2 > 0$, we choose δ such that $(\frac{1}{2\varepsilon} - r_2 - \delta) = C > 0$. Hence, (3.45) implies:

$$\begin{aligned} & \int_{\Omega} f_\varepsilon(\varphi_N) + \int_{\Omega} \int_{\Omega} M(x - y) \varphi_N(x) (1 - \varphi_N(y)) \\ & \geq C\|\varphi_N\|_H^2 + C. \quad (52) \end{aligned}$$

Combining the previous estimations, we find

$$\|\varphi_N(t)\|_H^2 + \int_0^t \int_{\Omega} \mu_\varepsilon(\varphi_N) |\nabla v_N|^2 \leq C + E \int_0^t \|\varphi_N\|_H^2 \quad (53)$$

Applying Gronwall's Lemma to (53) yields:

$$\|\varphi_N\|_{L^\infty(0, T, H)} \leq C, \quad (54)$$

and

$$\left\| \mu_\varepsilon^{\frac{1}{2}}(\varphi_N) |\nabla v_N| \right\|_{L^2(0, T, H)} \leq C. \quad (55)$$

Moreover, we have that

$$\begin{aligned} \int_{\Omega} \mu_{\varepsilon}(\varphi_N) |\nabla v_N|^2 &\geq \varepsilon \int_{\Omega} |\nabla f'_{\varepsilon}(\varphi_N) + \nabla M * (1 - 2\varphi_N)|^2 \\ &\geq \varepsilon \int_{\Omega} \left| \frac{(1 + 2b_{\varepsilon})\nabla \varphi_N}{\mu_{\varepsilon}(\varphi_N)} + \nabla M * (1 - 2\varphi_N) \right|^2 \\ &\geq p \int_{\Omega} |\nabla \varphi_N|^2 + C \int_{\Omega} |\nabla M * (1 - 2\varphi_N)|^2 \\ &\quad + C \int_{\Omega} |\nabla \varphi_N| |\nabla M * (1 - 2\varphi_N)| \\ &\geq p \int_{\Omega} |\nabla \varphi_N|^2 - Cr_2^2 \|1 - 2\varphi_N\|_H^2 \\ &\quad - \delta \int_{\Omega} |\nabla \varphi_N|^2 - \frac{r_2^2}{4\delta} \|1 - 2\varphi_N\|_H^2 \\ &\geq (p - \delta) \|\nabla \varphi_N\|_H^2 - (C + R_{\delta}) r_2^2 (\|\varphi_N\|_H^2 + 1). \end{aligned}$$

Here, π denotes a positive constant that depends solely on ε . Taking δ sufficiently small, we obtain:

$$\int_{\Omega} \mu_{\varepsilon}(\varphi_N) |\nabla v_N|^2 \geq p \|\varphi_N\|_H^2 - C \|\varphi_N\|_H^2 - C \quad (\clubsuit)$$

Substituting $(\clubsuit)$ into (3.47), we obtain:

$$\|\varphi_N\|_{L^2(0,T,Y)} \leq E, \quad (56)$$

and

$$\|\nabla v_N\|_{L^2(0,T,H)} \leq E. \quad (57)$$

Furthermore,

$$\begin{aligned} |\overline{v_N}| &= \left| \frac{1}{|\Omega|} \int_{\Omega} v_N \right| \\ &\leq E \left| \int_{\Omega} \frac{(1 + 2b_{\varepsilon})\nabla \varphi_N}{\mu_{\varepsilon}(\varphi_N)} \right| + E \int_{\Omega} |\nabla M * (1 - 2\varphi_N)| \\ &\leq E \left(\frac{R}{2} \|\nabla \varphi_N\|_H^2 + \frac{R}{2} \right) + \frac{Er_2^2}{2} \|1 - 2\varphi_N\|_H^2 + \frac{Er_2^2}{2} \\ &\leq E \|\varphi_N\|_H^2 + E + E \|\varphi_N\|_H^2 + E \\ &\leq E \|\varphi_N\|_H^2 + E \leq E \end{aligned} \quad (58)$$

which yields from Poincare-Wirtinger inequality and (58) that

$$\|v_N\|_{L^2(0,T,Y)} \simeq \|v_N\|_{L^2(0,T,V)} + |\overline{v_N}| \leq E + E \leq E. \quad (59)$$

Next, to obtain an a priori estimate for ϕ'_N , we take $\varphi = \phi'_N$ as a test function in (114), and employ (38), (57) and (59), we obtain the equivalent formulation:

$$\begin{aligned} ((\varphi'_N, \varphi'_N)) &= -((\mu_{\varepsilon}(\varphi_N) \nabla v_N, \nabla \varphi'_N)) \\ &\quad + ((\chi_{\Omega \setminus D}(x)(h - \varphi_N), \varphi'_N)) \\ &\leq E \|v_N\|_H \|\nabla \varphi_N\|_H + \|h - \varphi_N\|_H \|\varphi'_N\|_H \\ &\leq (E \|v_N\|_H + R \|\varphi_N\|_H + R) \|\varphi'_N\|_H \\ &\leq (C + C + R) \|\varphi'_N\|_H \leq C \|\varphi'_N\|_H. \end{aligned}$$

Thus,

$$\|\varphi'_N\|_{L^2(0,T,Y^*)} \leq C \quad (60)$$

As a result, from (57) and (57), we observe that φ_N , φ'_N , v_N and $f'_{\varepsilon}(\varphi_N)$ are bounded sequences in $L^2(\Omega)$ and $L^{\infty}(\Omega)$. Then there exists (not relabeled) subsequences of each such that:

$$\varphi_N \rightharpoonup \varphi \text{ weakly in } L^2(0, T, Y), \quad (61)$$

$$\varphi_N \rightharpoonup \varphi \text{ weakly in } L^{\infty}(0, T, H), \quad (62)$$

$$\varphi'_N \rightharpoonup \varphi' \text{ weakly in } L^2(0, T, Y^*), \quad (63)$$

$$f'_{\varepsilon}(\varphi_N) \rightharpoonup f^*_{\varepsilon} \text{ weakly in } L^2(0, T, Y), \quad (64)$$

and

$$v_N \rightharpoonup v \text{ weakly in } L^2(0, T, Y). \quad (65)$$

Remark 3.5 Let $(1 < p < \infty)$. Consider a Hilbert space $Y \subseteq H \subseteq Y^*$ and a sequence $\{\varphi_N\}$ where $\varphi_N : [0, T] \rightarrow Y$. Suppose that

$$\|\varphi_N\|_{L^2(0,T,Y)} \leq E, \quad \|\varphi'_N\|_{L^p(0,T,Y^*)} \leq E,$$

where $E > 0$ is independent of N . Then there $\exists$ a subsequence $\{\varphi_{N_i}\}$ such that

$$\varphi_{N_i} \rightharpoonup \varphi \text{ in } L^2(0, T, H).$$

Moreover, using Remark 3.5 and owing to (56) and (60), we get:

$$\varphi_N \rightarrow \varphi \text{ strongly in } L^2(0, T, H) \text{ and a.e. in } \psi. \quad (66)$$

Now (3.44) and (3.48) imply the following:

$$\|\chi_{\Omega \setminus D}(x)(h - \varphi_N)\|_{L^2(0,T,H)} \leq E. \quad (67)$$

Additionally, μ_{ε} and $\chi_{\Omega \setminus D}(x)(h - \varphi_N)$ are continuous functions. Then, by applying the Dominated Convergence Theorem, we get

$$\mu_{\varepsilon}(\varphi_N) \rightarrow \mu_{\varepsilon}(\varphi) \text{ a.e. in } \psi, \quad (68)$$

and

$$\chi_{\Omega \setminus D}(x)(h - \varphi_N) \rightarrow \chi_{\Omega \setminus D}(x)(h - \varphi) \text{ in } L^2(0, T, H). \quad (69)$$

Hence, using (65) and (69), we reach:

$$\mu_{\varepsilon}(\varphi_N) v_N \rightarrow \mu_{\varepsilon}(\varphi) v \text{ weakly in } L^2(\psi). \quad (70)$$

In fact, suppose $\phi \in L^2(0, T, H)$. Then, for $\{j \in 1, \dots, l\}$ we have:

$$\int_0^T (\mu_{\varepsilon}(\varphi_N) \partial_j v_N, \phi) = \int_0^T (\partial_j v_N, \mu_{\varepsilon}(\varphi_N) \phi). \quad (71)$$

Also, (71) is equivalent to:

$$\begin{aligned} \int_0^T (\mu_\varepsilon(\varphi_N) \partial_j v_N, \phi) &= \int_0^T (\partial_j v_N, \mu_\varepsilon(\varphi) \phi) \\ &+ \int_0^T (\partial_j v_N - \partial_j v, \mu_\varepsilon(\varphi) \phi) \\ &+ \int_0^T (\partial_j v_N, (\mu_\varepsilon(\varphi_N) - \mu_\varepsilon(\varphi)) \phi). \end{aligned} \quad (72)$$

Again applying the Dominated Convergence Theorem, we find

$$\begin{aligned} &\left| \int_0^T (\partial_j v_N, (\mu_\varepsilon(\varphi_N) - \mu_\varepsilon(\varphi)) \phi) \right| \\ &\leq \|\partial_j v_N\|_{L^2(0,T,H)} \|(\mu_\varepsilon(\varphi_N) - \mu_\varepsilon(\varphi)) \phi\|_{L^2(0,T,H)} \\ &\leq E \|(\mu_\varepsilon(\varphi_N) - \mu_\varepsilon(\varphi)) \phi\|_{L^2(0,T,H)}. \end{aligned} \quad (\star)$$

Hence

$$\left| \int_0^T (\partial_j v_N - \partial_j v, \mu_\varepsilon(\varphi) \phi) \right| \quad (\star\star)$$

$$\leq \|\partial_j v_N - \partial_j v\|_{L^2(0,T,H)} \|\mu_\varepsilon(\varphi) \phi\|_{L^2(0,T,H)}.$$

Now, as $N \rightarrow \infty$, we have $(\star) \rightarrow 0$ and $(\star\star) \rightarrow 0$. Thus, we conclude that:

$$\int_0^T (\mu_\varepsilon(\varphi_N) \partial_j v_N, \phi) \longrightarrow 0 \text{ as } N \rightarrow \infty.$$

Lastly, using the fact that f'_ε is continuous, we find that $f'_\varepsilon(\varphi) = f_\varepsilon^*$. Taking the limit N to ∞ , and applying the previous convergence results, we conclude that φ is a solution of (9).

In addition, using Fatou Lemma, we reach:

$$\begin{aligned} &\left\| \mu_\varepsilon^{\frac{1}{2}}(\varphi) |\nabla v| \right\|_{L^2(0,T,L^2(\Omega))} \\ &\leq \liminf_{N \rightarrow \infty} \left\| \mu_\varepsilon^{\frac{1}{2}}(\varphi) |\nabla v| \right\|_{L^2(0,T,L^2(\Omega))} \leq C \end{aligned}$$

Thus, the proof of Lemma 3.4 is complete. $\square$
Now, we multiply (29) by the test function $\phi = \varphi_\varepsilon$ and, applying (A), (38) and Young's inequality, we obtain:

$$\begin{aligned} ((\varphi'_\varepsilon, \varphi_\varepsilon)) &= \frac{1}{2} \frac{d}{dt} \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 \\ &= - \int_\Omega \mu_\varepsilon \nabla \varphi_\varepsilon \nabla v_\varepsilon \\ &+ \int_{\Omega \setminus D} \varphi_\varepsilon \chi_{\Omega \setminus D} (h - \varphi_\varepsilon) \end{aligned}$$

$$\begin{aligned} &= - \int_\Omega \mu_\varepsilon \nabla \varphi_\varepsilon f'_\varepsilon(\varphi_\varepsilon) - \int_\Omega \mu_\varepsilon \nabla \varphi_\varepsilon \nabla \theta_\varepsilon \\ &+ \int_{\Omega \setminus D} \varphi_\varepsilon \chi_{\Omega \setminus D} (h - \varphi_\varepsilon) \\ &\leq - \int_\Omega |\nabla \varphi_\varepsilon|^2 + Cr_2 \|\nabla \varphi_\varepsilon\|_{L^2(\Omega)} \|1 - 2\varphi_\varepsilon\|_{H^1(\Omega)} \\ &+ \frac{1}{2} \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + \frac{1}{2} \|h\|_{L^2(\Omega)}^2 - \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 \\ &\leq - \int_\Omega |\nabla \varphi_\varepsilon|^2 + Cr_2 (\delta \|\nabla \varphi_\varepsilon\|_{L^2(\Omega)}^2) \\ &+ R_\delta \|1 - 2\varphi_\varepsilon\|_{H^1(\Omega)}^2 + E \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + E \\ &\leq - \int_\Omega |\nabla \varphi_\varepsilon|^2 + Cr_2 (\delta \|\nabla \varphi_\varepsilon\|_{L^2(\Omega)}^2) \\ &+ R_\delta \|1 - 2\varphi_\varepsilon\|_{H^1(\Omega)}^2 + E \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + E \\ &\leq (\delta - 1) \int_\Omega |\nabla \varphi_\varepsilon|^2 + R_\delta \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + R_\delta \forall \delta > 0, \end{aligned}$$

where R_δ is a constant depending on δ . This inequality is equivalent to:

$$\frac{1}{2} \frac{d}{dt} \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + (1 - \delta) \|\nabla \varphi_\varepsilon\|_{L^2(\Omega)}^2 \leq R_\delta \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + R_\delta.$$

Let $\delta < 1$ be fixed. Then:

$$\frac{1}{2} \frac{d}{dt} \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + \|\nabla \varphi_\varepsilon\|_{L^2(\Omega)}^2 \leq R \|\varphi_\varepsilon\|_{L^2(\Omega)}^2 + R.$$

Applying Grownall's Lemma, we get:

$$\|\varphi_\varepsilon\|_{L^\infty(0,T,L^2(\Omega))} + \|\nabla \varphi_\varepsilon\|_{L^2(0,T,L^2(\Omega))} \leq R, \quad (73)$$

from which we can deduce that:

$$\|\varphi_\varepsilon\|_{L^2(0,T,H^1(\Omega))} < R. \quad (74)$$

Likewise, choosing the test function $\phi = v_\varepsilon$ in (29), we obtain:

$$\begin{aligned} &((\varphi'_\varepsilon, v_\varepsilon)) + ((m_\varepsilon \nabla v_\varepsilon, \nabla v_\varepsilon)) \\ &= \frac{d}{dt} \left\{ \int_\Omega f_\varepsilon(\varphi_\varepsilon) + \int_\Omega [M * (1 - \varphi_\varepsilon)] \varphi_\varepsilon \right\} \\ &+ \int_\Omega \mu_\varepsilon |\nabla v_\varepsilon|^2 = \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x) (h - \varphi_\varepsilon) v_\varepsilon \end{aligned} \quad (75)$$

Owing to (A) and (74), we find

$$\begin{aligned} \left| \int_\Omega [M * (1 - \varphi_\varepsilon)] \varphi_\varepsilon \right| &\leq \|\varphi_\varepsilon\|_{L^2(\Omega)} \|M * (1 - \varphi_\varepsilon)\|_{L^2(\Omega)} \\ &\leq R \|1 - 2\varphi_\varepsilon\|_{L^2(\Omega)} \\ &\leq R(|\Omega| - \|\varphi_\varepsilon\|_{L^2(\Omega)}) \leq R, \end{aligned} \quad (76)$$

and

$$\begin{aligned}
 & \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)v_\varepsilon \\
 &= \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)f'_\varepsilon + \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)\theta_\varepsilon \\
 &\leq \int_{\varphi_\varepsilon \leq 0} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)f'_\varepsilon + \int_{\varphi_\varepsilon \geq 0} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)f'_\varepsilon \\
 &+ \int_{\varphi_\varepsilon \in (0,1)} \chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon)f'_\varepsilon \\
 &+ \int_{\Omega} (h - \varphi_\varepsilon)[M * (1 - 2\varphi_\varepsilon)] \\
 &\leq \int_{\varphi_\varepsilon \in (0,1)} (h - \varphi_\varepsilon) \ln \left(\frac{\varphi_\varepsilon + b_\varepsilon}{1 - \varphi_\varepsilon + b_\varepsilon} \right) + Cr_2 \|1 - 2\varphi_\varepsilon\|_{L^2(\Omega)} \\
 &\leq \int_{\varphi_\varepsilon \in (0,1)} (h - \varphi_\varepsilon) \ln(\varphi_\varepsilon + b_\varepsilon) \\
 &- \int_{\varphi_\varepsilon \in (0,1)} (h - \varphi_\varepsilon) \ln(1 - \varphi_\varepsilon + b_\varepsilon) + R \\
 &\leq (R - \|h\|_{L^2(\Omega)})E - (R - \|h\|_{L^2(\Omega)})C + R \\
 &\leq R.
 \end{aligned} \tag{77}$$

Thanks to (76) and (77), we have that

$$\left\| \mu_\varepsilon^{\frac{1}{2}} |\nabla v_\varepsilon| \right\|_{L^2(\psi)} \leq R, \tag{78}$$

and

$$\left| \int_{\Omega} f_\varepsilon(\varphi_\varepsilon) \right| \leq R. \tag{79}$$

Similarly, for φ_ε , let $\phi = \varphi'_\varepsilon$ be the test function in (29). Then, using (2.30), and (2.67), we obtain:

$$\|\varphi'_\varepsilon\|_{L^2(0,T,(H^1(\Omega))^*)} \leq R. \tag{80}$$

Hence, owing to Remark 2 and (77), there exists a (not relabeled) subsequence $\varphi \in H^1(0, T, H^1(\Omega), (H^1(\Omega))^*) \cap L^\infty(0, T, L^2(\Omega))$ such that:

$$\varphi_\varepsilon \rightharpoonup \varphi \text{ weakly in } L^2(0, T, H^1(\Omega)), \tag{81}$$

$$\varphi_\varepsilon \rightharpoonup \varphi \text{ weakly in } L^\infty(0, T, L^2(\Omega)), \tag{82}$$

$$\varphi'_\varepsilon \rightharpoonup \varphi' \text{ weakly in } L^2(0, T, (H^1(\Omega))^*), \tag{83}$$

$$\chi_{\Omega \setminus D}(x)(h - \varphi_\varepsilon) \rightharpoonup \chi_{\Omega \setminus D}(x)(h - \varphi) \tag{84}$$

pointwise a.e in ψ , and

$$\varphi_\varepsilon \rightharpoonup \varphi \text{ strongly in } L^2(0, T, L^2(\Omega)) \text{ a.e in } \psi. \tag{85}$$

In addition, we have

$$\theta_\varepsilon \rightarrow \theta = M * (1 - 2\varphi) \text{ in } L^2(0, T, H^1(\Omega)) \tag{86}$$

and

$$\mu_\varepsilon(\varphi_\varepsilon) \rightarrow \mu(\varphi) \text{ a.e. in } \psi. \tag{87}$$

Thus, by combining (86) and (87), we reach:

$$\mu_\varepsilon(\varphi_\varepsilon) \nabla \theta_\varepsilon \rightarrow \mu(\varphi) \nabla \theta \text{ in } L^2(\psi). \tag{88}$$

On the other hand, applying the Dominated Convergence Theorem, we find

$$\begin{aligned}
 & \|\mu_\varepsilon(\varphi_\varepsilon) \nabla \theta_\varepsilon - \mu(\varphi) \nabla \theta\|_{L^2(\Omega)} \\
 &\leq \|(\mu_\varepsilon(\varphi_\varepsilon) - \mu(\varphi)) \nabla \theta_\varepsilon\|_{L^2(\Omega)} + \|\mu(\varphi) (\nabla \theta_\varepsilon - \nabla \theta)\|_{L^2(\Omega)} \\
 &\leq \|\mu_\varepsilon(\varphi_\varepsilon) - \mu(\varphi)\|_{L^2(\Omega)} \|\nabla \theta_\varepsilon\|_{L^2(\Omega)} \\
 &+ \|\mu(\varphi)\|_{L^2(\Omega)} \|\nabla \theta_\varepsilon - \nabla \theta\|_{L^2(\Omega)} \rightarrow 0 \text{ as } \varepsilon \rightarrow 0.
 \end{aligned}$$

Therefore, taking the limit $\varepsilon \rightarrow 0$, we deduce that φ is a solution of (9)-(13).

Besides, we have that

$$\mu_\varepsilon(m) = \varepsilon \quad \forall m < 0 \text{ and } \forall m > 1$$

and

$$f''_\varepsilon(m) = \frac{1 + 2b_\varepsilon}{\varepsilon} \quad \forall m < 0 \text{ and } \forall m > 1.$$

Hence, for $m > 1$, we reach:

$$f'_\varepsilon(m) \geq \frac{1 + 2b_\varepsilon}{\varepsilon} (m - 1) + f'_\varepsilon(m) \geq \frac{1 + 2b_\varepsilon}{\varepsilon} (m - 1)$$

and

$$f_\varepsilon(m) \geq \frac{1 + 2b_\varepsilon}{2\varepsilon} (m - 1)^2 + f_\varepsilon(1).$$

Similarly for $m < 0$, we have:

$$f_\varepsilon(m) \geq \frac{1 + 2b_\varepsilon}{2\varepsilon} m^2 + f_\varepsilon(0).$$

Now, for $m = \varphi_\varepsilon$, we have

$$\begin{aligned}
 & \int_{\varphi_\varepsilon > 1} (\varphi_\varepsilon - 1)^2 \leq \frac{2\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \int_{\varphi_\varepsilon > 1} |f_\varepsilon(\varphi_\varepsilon)| \\
 &- \frac{2\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \int_{\varphi_\varepsilon > 1} |f_\varepsilon(1)| \\
 &\leq \frac{\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \left(R - 2 \int_{\varphi_\varepsilon > 1} |f_\varepsilon(1)| \right)
 \end{aligned} \tag{89}$$

and

$$\begin{aligned}
 & \int_{\varphi_\varepsilon < 0} \varphi_\varepsilon^2 \\
 &\leq \frac{2\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \int_{\varphi_\varepsilon < 0} |f_\varepsilon(\varphi_\varepsilon)| - \frac{2\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \int_{\varphi_\varepsilon < 0} |f_\varepsilon(1)| \\
 &\leq \frac{\mu_\varepsilon(1)}{1 + 2b_\varepsilon} \left(R - 2 \int_{\varphi_\varepsilon < 0} |f_\varepsilon(1)| \right).
 \end{aligned} \tag{90}$$

Choosing $\frac{\mu_\varepsilon(1)}{1+2b_\varepsilon} = o(1)$ and $f_\varepsilon = o(1)$ as $\varepsilon \rightarrow 0$ in both (89) and (90), we obtain:

$$\int_{\varphi>1} (\varphi - 1)^2 = 0$$

and

$$\int_{\varphi<0} \varphi^2 = 0.$$

Therefore,

$$0 \leq \varphi \leq 1 \text{ a.e in } \psi.$$

In conclusion, Theorem 3.2 is proved.

3.1 Further regularity

Theorem 3.6 Suppose the assumptions of Theorem 3.2 are satisfied. Then, there exists a weak solution to (4)-(8) as in Definition 3.1. Additionally, φ_0 satisfies:

$$\varphi_0 \in H^2(\Omega) \quad (91)$$

and

$$n \cdot (\nabla \varphi_0 + \mu(\varphi_0) \nabla M * (1 - 2\varphi_0)) = 0 \text{ on } \partial\Omega. \quad (92)$$

Then, $\varphi \in L^\infty(0, T, H^2(\Omega))$.

The four lemmas, along with there proofs, are as follows:

Lemma 3.7 Assume the conditions of Theorem 3.1 are satisfied. Then, there exists φ a solution of (4)-(8) as defined in Definition 3.1 such that

$$\varphi_t \in L^\infty(0, T, (H^1(\Omega))^* \cap L^2(0, T, L^2(\Omega))). \quad (93)$$

Proof: Firstly, we have $\varphi_t(0) \in (H^1(\Omega))^*$. Consequently, we obtain:

$$\bar{\varphi}_t = \int_{\Omega} \chi_{\Omega \setminus D}(x)(\varphi - h)dx \text{ for a.e. } t \in [0, T]. \quad (94)$$

However, since Ω is bounded, we have:

$$\bar{\varphi}_t \in L^\infty(0, T) \text{ and } \|\bar{\varphi}_t\|_{L^\infty(0, T)} \leq E. \quad (95)$$

Denote $(H_0^1(\Omega))^*$ by $H_0^{-1}(\Omega)$. Then, it remains to show that:

$$\Phi = \varphi_t - \bar{\varphi}_t \in L^\infty(0, T, H_0^{-1}(\Omega)) \cap L^2(0, T, L^2(\Omega)).$$

Consider the realization of the Laplacian with the boundary conditions $\Delta_N : H_0^1(\Omega) \rightarrow H_0^{-1}(\Omega)$. In other words, the proof will proceed

by approximating the time derivative with Δ_N . Now, choosing the test function $\phi = \Delta_N^{-1}\Phi_t$ and noting that $\bar{\varphi}_t$ has a zero mean, the variational formulation of (4) becomes:

$$((\Phi_t, \Delta_N^{-1}\Phi_t))$$

$$= -((\mu(\Phi)\nabla v, \nabla \Delta_N^{-1}\Phi_t)) + ((\chi_{\Omega \setminus D}(x)(\varphi - h), \Delta_N^{-1}\Phi_t)).$$

Differentiating each term in the variational formula above with respect to t , we reach the following computations:

$$\begin{aligned} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 &= 2((\Phi_{tt}, \Phi_t))_{H_0^{-1}(\Omega)} \\ &= 2((\nabla \Delta_N^{-1}\Phi_{tt}, \nabla \Delta_N^{-1}\Phi_t)) = -2((\Phi_{tt}, \Delta_N^{-1}\Phi_t))_{L^2(\Omega)}, \end{aligned}$$

which then implies:

$$\frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + ((\Phi_{tt}, \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} = 0. \quad (96)$$

Moreover, we have

$$\begin{aligned} &-((\mu(\Phi)(\nabla v)_t, \nabla \Delta_N^{-1}\Phi_t)) \\ &= -(((\mu \nabla \theta)_t, \nabla \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} - ((\nabla \Phi_t, \nabla \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} \\ &= -(((\mu \nabla \theta)_t, \nabla \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} + \|\Phi_t\|_{L^2(\Omega)}^2. \end{aligned}$$

From this, and using the identity

$$-((\mu(\nabla v)_t, \nabla \Delta_N^{-1}\Phi_t)) = ((\nabla(\mu \nabla v)_t, \Delta_N^{-1}\Phi_t)),$$

we obtain:

$$\begin{aligned} &((\nabla(\mu \nabla v)_t, \Delta_N^{-1}\Phi_t)) + (((\mu \nabla \theta)_t, \nabla \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} \\ &- \|\Phi_t\|_{L^2(\Omega)}^2 = 0. \end{aligned} \quad (97)$$

Thus, combining (96) and (97), we get:

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + ((\Phi_{tt}, \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} \\ &= -\|\Phi_t\|_{L^2(\Omega)}^2 + ((\nabla(\mu \nabla v)_t, \Delta_N^{-1}\Phi_t)) \\ &\quad + (((\mu \nabla \theta)_t, \nabla \Delta_N^{-1}\Phi_t))_{L^2(\Omega)}. \end{aligned}$$

Hence,

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + \|\Phi_t\|_{L^2(\Omega)}^2 \\ &= -((\Phi_{tt}, \Delta_N^{-1}\Phi_t))_{L^2(\Omega)} + (((\mu \nabla \theta)_t, \nabla \Delta_N^{-1}\Phi_t)) \\ &\quad + ((\nabla(\nabla \varphi + \mu \nabla \theta)_t, \Delta_N^{-1}\Phi_t))_{L^2(\Omega)}. \end{aligned} \quad (98)$$

Furthermore, differentiating (4) with respect to t , we get:

$$\varphi_{tt} = \nabla \varphi_t + (\mu \nabla \theta)_t + \chi_{\Omega \setminus D}(x)(\varphi_t - h_t) \quad (99)$$

and

$$\bar{\varphi}_{tt} = \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t)dx. \quad (100)$$

Then, Subtracting (99) and (100), we deduce the following equation:

$$\begin{aligned} \Phi_{tt} &= \varphi_{tt} - \bar{\varphi}_{tt} \\ &= \nabla \varphi_t + (\mu \nabla \theta)_t + \chi_{\Omega \setminus D}(x)(\varphi_t - h_t) \\ &\quad + \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \end{aligned}$$

that is equivalent to:

$$\begin{aligned} &-((\Phi_{tt}, \Delta_N^{-1} \phi_t))_{L^2(\Omega)} + ((\nabla(\nabla \theta + \mu \nabla \theta)_t, \Delta_N^{-1} \Phi_t)) \\ &= -((\chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t)) \\ &\quad + \left(\left(\int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t \right) \right). \end{aligned}$$

As a result, we deduce:

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + \|\Phi_t\|_{L^2(\Omega)}^2 &= (((\mu \nabla \theta)_t, \nabla \Delta_N^{-1} \Phi_t)) \\ &\quad - ((\chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t)) \\ &\quad + \left(\left(\int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t \right) \right). \end{aligned} \quad (101)$$

It then follows from Young's inequality that

$$\begin{aligned} &(((\mu \nabla \theta)_t, \nabla \Delta_N^{-1} \Phi_t)) \\ &\leq \|\mu_t \nabla \theta + \mu \nabla \theta_t\|_{L^2(\Omega)} \|\nabla \Delta_N^{-1} \Phi_t\|_{L^2(\Omega)} \\ &\leq r_\infty \|\varphi_t\|_{L^2(\Omega)} \|\Phi_t\|_{H_0^{-1}(\Omega)} + r_2 \|\varphi_t\|_{L^2(\Omega)} \|\Phi_t\|_{H_0^{-1}(\Omega)} \\ &\leq R \|\varphi_t\|_{L^2(\Omega)} \|\Phi_t\|_{H_0^{-1}(\Omega)} \\ &\leq R(1 + \|\Phi_t\|_{L^2(\Omega)}) \|\Phi_t\|_{H_0^{-1}(\Omega)} \\ &\leq R \|\Phi_t\|_{H_0^{-1}(\Omega)} + R \left(\frac{1}{2} \|\Phi_t\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 \right) \\ &\leq \frac{1}{2} \|\Phi_t\|_{L^2(\Omega)}^2 + E \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + R \|\Phi_t\|_{H_0^{-1}(\Omega)}. \end{aligned} \quad (102)$$

Using the fact that $\Delta_N^{-1} \Phi_t \in H_0^1(\Omega)$ and Poincare's inequality, we find

$$\begin{aligned} &((\chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t)) \\ &= \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t) \Delta_N^{-1} \Phi_t \\ &\leq \|\varphi_t - h_t\|_{L^2(\Omega)} \|\Delta_N^{-1} \Phi_t\|_{L^2(\Omega)} \\ &\leq (1 + \|\Phi_t\|_{L^2(\Omega)} - \|h_t\|_{L^2(\Omega)}) \|\Delta_N^{-1} \Phi_t\|_{L^2(\Omega)} \\ &\leq \frac{1}{4} \|\Phi_t\|_{L^2(\Omega)}^2 + R \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E \|\Phi_t\|_{H_0^{-1}(\Omega)} \end{aligned} \quad (103)$$

and

$$\begin{aligned} &\left(\left(\int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Delta_N^{-1} \Phi_t \right) \right) \\ &\leq \|\varphi_t - h_t\|_{L^2(\Omega)} \|\Delta_N^{-1} \Phi_t\|_{L^2(\Omega)} \\ &\leq (|\Omega| + \|\Phi_t\|_{L^2(\Omega)} - \|h_t\|_{L^2(\Omega)}) \|\Delta_N^{-1} \Phi_t\|_{L^2(\Omega)} \\ &\leq R|\Omega|^2 \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E|\Omega| \|\Phi_t\|_{H_0^{-1}(\Omega)} + \frac{1}{8} \|\Phi_t\|_{L^2(\Omega)}^2. \end{aligned} \quad (104)$$

Thus, by combining (102), (103), (104), and applying Young's inequality in (101), we obtain:

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + \frac{1}{8} \|\Phi_t\|_{L^2(\Omega)}^2 \\ &\leq R \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E \|\Phi_t\|_{H_0^{-1}(\Omega)} \\ &\leq R \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E \left(\frac{1}{2} \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + \frac{1}{2} \right) \\ &\leq E \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E \end{aligned}$$

Integrating over $(0, t)$ gives:

$$\begin{aligned} &\frac{1}{2} \|\Phi_t(t)\|_{H_0^{-1}(\Omega)}^2 + \frac{1}{8} \int_0^t \|\Phi_t\|_{L^2(\Omega)}^2 \\ &\leq E \int_0^t \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E + \frac{1}{2} \|\Phi_0\|_{H_0^{-1}(\Omega)} \end{aligned}$$

that also implies:

$$\|\Phi_t(t)\|_{H_0^{-1}(\Omega)}^2 + \frac{1}{4} \int_0^t \|\Phi_t\|_{L^2(\Omega)}^2 \leq E \int_0^t \|\Phi_t\|_{H_0^{-1}(\Omega)}^2 + E.$$

By now applying Gronwall's Lemma, we have:

$$\|\Phi_t\|_{L^\infty(0,T,(H^1(\Omega))^*)} + \|\Phi_t\|_{L^2(0,T,L^2(\Omega))} \leq E.$$

Finally, by using (3.5), we obtain the estimate:

$$\|\varphi_t\|_{L^\infty(0,T,(H^1(\Omega))^*)} + \|\varphi_t\|_{L^2(0,T,L^2(\Omega))} \leq E. \quad \square$$

Lemma 3.8 *Given that the conditions of Theorem 3.1 are satisfied, there exists φ a solution to the system (4)-(8) as defined in Definition 2.1, such that*

$$\varphi_t \in L^\infty(0, T, L^2(\Omega)) \cap L^2(0, T, H^1(\Omega)). \quad (105)$$

Proof: We select the test function $\phi = \Phi_t$ and differentiate (4) with respect to time. This gives:

$$\begin{aligned}
 \frac{1}{2} \frac{d}{dt} \|\Phi_t\|_{L^2(\Omega)}^2 &= -((\mu_t \nabla \theta + \mu \nabla \theta_t + \nabla \Phi_t, \nabla \Phi_t))_{L^2(\Omega)} \\
 &+ ((\chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Phi_t))_{L^2(\Omega)} \\
 &- \left(\left(\int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x)(\varphi_t - h_t), \Phi_t \right) \right)_{L^2(\Omega)} \\
 &\leq -R \|\nabla \Phi_t\|_{L^2(\Omega)}^2 - \|\Phi_t\|_{L^2(\Omega)} \|\nabla \Phi_t\|_{L^2(\Omega)} \\
 &+ R \|\Phi_t\|_{L^2(\Omega)} + \|\Phi_t\|_{L^2(\Omega)}^2 - E \|\Phi_t\|_{L^2(\Omega)} + \|\Phi_t\|_{L^2(\Omega)}^2 \\
 &\leq -R \left(\frac{1}{2} \|\Phi_t\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\nabla \Phi_t\|_{L^2(\Omega)}^2 \right) - R \|\nabla \Phi_t\|_{L^2(\Omega)}^2 \\
 &+ R \|\Phi_t\|_{L^2(\Omega)}^2 + E \|\Phi_t\|_{L^2(\Omega)} \\
 &\leq R \|\Phi_t\|_{L^2(\Omega)} + E \|\Phi_t\|_{L^2(\Omega)}^2 - \frac{1}{2} \|\nabla \Phi_t\|_{L^2(\Omega)}^2.
 \end{aligned} \tag{106}$$

Then, by integrating (106) over $(0, t)$, we infer:

$$\begin{aligned}
 \frac{1}{2} \frac{d}{dt} \|\Phi_t(t)\|_{L^2(\Omega)}^2 + \frac{1}{2} \int_0^t \|\nabla \Phi_t\|_{L^2(\Omega)}^2 \\
 \leq R \int_0^t \|\Phi_t\|_{L^2(\Omega)} + E \int_0^t \|\Phi_t\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\Phi_t(0)\|_{L^2(\Omega)}^2.
 \end{aligned}$$

Additionally, we observe that $\|\Phi_t(0)\|_{L^2(\Omega)}$ is bounded. Then, using Young's inequality, we reach:

$$\frac{d}{dt} \|\Phi_t(t)\|_{L^2(\Omega)}^2 + \int_0^t \|\nabla \Phi_t\|_{L^2(\Omega)}^2 \leq R \int_0^t \|\Phi_t\|_{L^2(\Omega)}^2 + E.$$

Moreover, by applying Gronwall's Lemma, we obtain:

$$\|\Phi_t\|_{L^\infty(0, T, L^2(\Omega))} + \|\nabla \Phi_t\|_{L^2(0, T, H^1(\Omega))} \leq E.$$

Then, using the fact that $\nabla \bar{\varphi}_t = 0$, we get

$$\|\varphi_t\|_{L^\infty(0, T, L^2(\Omega))} + \|\varphi_t\|_{L^2(0, T, H^1(\Omega))} \leq E.$$

Thus, (105) follows. $\square$

Lemma 3.9 Suppose the conditions of Theorem 3.1 are satisfied. Then, the solution φ to the system (4)-(8) as defined in Definition 2.1, satisfies:

$$\nabla \varphi \in L^\infty(0, T, L^2(\Omega)). \tag{107}$$

Proof: We deduce that $\nabla \varphi \in H^1(0, T, L^2(\Omega))$, owing to (105). Therefore, we obtain (111). $\square$

Lemma 3.10 Assume that the conditions of Theorem 3.1 are satisfied. Then, there exists φ a solution to the system (4)-(8), as defined in Definition 2.1, such that:

$$\varphi \in L^\infty(0, T, H^2(\Omega)). \tag{108}$$

Proof: Owing to Lemma 3.2 and Lemma 3.3, we have $\varphi_t \in L^\infty(0, T, L^2(\Omega))$, $(1 - 2\varphi)\nabla \varphi \nabla \theta \in L^\infty(0, T, L^2(\Omega))$ and $\mu \Delta \theta \in L^\infty(0, T, L^2(\Omega))$. From this, we can also conclude that $\chi_{\Omega \setminus D}(x)(\varphi - h) \in L^\infty(0, T, L^2(\Omega))$. Hence,

$$((\Delta \varphi, \phi)) = ((\zeta, \phi)) \quad \forall \phi \in H^1(\Omega) \text{ a.e. } t \in (0, T), \tag{109}$$

such that

$$\zeta = \varphi_t + (1 - 2\varphi)\nabla \varphi \nabla \theta + \mu \Delta \theta + \chi_{\Omega \setminus D}(x)(\varphi - h) \tag{110}$$

in $L^\infty(0, T, L^2(\Omega))$. We thus infer that

$$\varphi \in L^\infty(0, T, H^1(\Omega)) \cap L^\infty(\psi)$$

$$\theta \in L^\infty(0, T, H^2(\Omega)) \cap L^\infty(0, T, W^{1, \infty}(\Omega))$$

$$\nabla \theta \in L^\infty(0, T, H^1(\Omega)) \cap L^\infty(\psi).$$

Moreover, since $\partial\Omega$ is locally Lipschitz, there exists a unit normal vector n on $\partial\Omega$ such that $n \in L^\infty(\partial\Omega)$. Thus,

$$\frac{\partial \theta}{\partial n} \in L^\infty\left(0, T, H^{\frac{1}{2}}(\partial\Omega)\right) \cap L^\infty(0, T, L^\infty(\partial\Omega)). \tag{111}$$

By Lemma 3.3, we have $\nabla \mu(\varphi) = (1 - 2\varphi)\nabla \varphi \in L^\infty(0, T, L^2(\Omega))$, from which we deduce:

$$\mu(\varphi) \in L^\infty\left(0, T, H^{\frac{1}{2}}(\Omega)\right). \tag{112}$$

Given that $0 \leq \varphi \leq 1$ and $0 \leq \mu(m) \leq 1$ for all $m \in [0, 1]$, and noting that this can be extended for all $m \in \mathbb{R}$ such that $0 \leq \mu(m) \leq 1$, we conclude:

$$\mu(\varphi) \in L^\infty(0, T, L^\infty(\Omega)). \tag{113}$$

Now, from (111), (112) and (113), we infer:

$$\mu(\varphi) \frac{\partial \theta}{\partial n} \in L^\infty\left(0, T, H^{\frac{1}{2}}(\partial\Omega)\right). \tag{114}$$

It then follows that:

$$\frac{\partial \varphi}{\partial n} = n \cdot \mu \nabla \theta = \mu(\varphi) \frac{\partial \theta}{\partial n} \text{ a.e on } \partial\Omega.$$

Then, by using (114), we deduce:

$$\frac{\partial \varphi}{\partial n} \in L^\infty\left(0, T, H^{\frac{1}{2}}(\partial\Omega)\right). \tag{115}$$

Remark 3.11 Elliptic Regularity Theorem: Let $\Omega \subset \mathbb{R}^l$, where $l \in \mathbb{N}$. Denote by n the outer unit normal on $\partial\Omega$. Suppose $\zeta \in L^2(\Omega)$ and $\kappa \in$

$H^{\frac{1}{2}}(\partial\Omega)$. Then, $\rho \in H^2(\Omega)$, if there exists $\rho \in H^1(\Omega)$, such that ρ is a weak solution satisfying:

$$\begin{cases} \Delta\rho = \zeta \text{ in } \Omega, \\ \frac{\partial\rho}{\partial n} = \kappa \text{ on } \partial\Omega. \end{cases}$$

Furthermore, there exists $E > 0$, independent of ζ and κ , such that

$$\|\rho\|_{H^2(\Omega)} \leq E \left(\|\rho\|_{L^2(\Omega)} + \|\zeta\|_{L^2(\Omega)} + \|\kappa\|_{H^{\frac{1}{2}}(\partial\Omega)} \right),$$

Lastly, by using Remark 3, we obtain:

$$\|\varphi\|_{H^2(\Omega)} \leq E \left(\|\varphi\|_{L^2(\Omega)} + \|\zeta\|_{L^2(\Omega)} + \left\| \frac{\partial\varphi}{\partial n} \right\|_{H^{\frac{1}{2}}(\partial\Omega)} \right)$$

for almost every $t \in [0, T]$. Considering (120), (115), and the result $\varphi \in L^\infty(0, T, L^2(\Omega))$, we arrive at (108). $\square$

3.2 The Separation Property

This part focuses on the separation of the solution φ of the system (4)-(8) from singularities. First, we show that $\bar{\varphi}(t)$ lies between 0 and 1. After obtaining some a priori estimates, we prove the existence of a positive constant q such that $q \leq \varphi(x, t) \leq 1 - q$ a.e. $x \in \Omega$, $t \in [T_0, T]$. Thus, we have following theorem:

Theorem 3.12 *Given $l \leq 3$, and assuming the conditions of Theorem 3.1 are verified, then for all $T_0 \in (0, T)$, there exists a constant $q > 0$ such that:*

$$q \leq \varphi(x, t) \leq 1 - q \text{ for almost every } x \in \Omega, t \in [T_0, T]. \quad (116)$$

Note that if

$$\exists \tilde{q} > 0 \text{ such that } \tilde{q} \leq \varphi_0 \leq 1 - \tilde{q}, \quad (117)$$

then $T_0 = 0$.

We begin by proving that $0 < \bar{\varphi}(t) < 1$ for all $t \in [0, T]$ in Lemmas 4.1 and 4.2.

Remark 3.13 *If $0 \leq \varphi \leq 1$, then φ separates 0 and 1 in the following manner:*

$$0 < \bar{\varphi}(t) = \frac{1}{|\Omega|} \int_{\Omega} \varphi(x, t) dx < 1 \quad \forall t \in [0, T].$$

Lemma 3.14 *Assume that φ is a weak solution of the system (2.1)-(2.5) as defined in Definition 2.1 and that (B3) is satisfied. Then, there exist constants $d_0 > 0$ and $e_0 > 0$, independent of t , such that:*

$$|\Omega_1^t| \geq e_0 > 0, \quad (118)$$

where $\Omega_1^t = \{x' \in \Omega; \varphi(x', t) \geq d_0\}$.

Proof: For simplicity, we assume $|\Omega| = 1$. We have

$$\bar{\varphi}(t) - \bar{\varphi}_0 = \int_0^t \frac{d}{dm} \bar{\varphi}(m) dm = \int_0^t ((\varphi', 1))$$

$$= \int_0^t \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x) (\varphi - h) = \int_0^t ((\chi_{\Omega \setminus D}(x) (\varphi - h), 1)).$$

Consequently, we deduce:

$$\bar{\varphi}(t) = \bar{\varphi}_0 + \int_0^t \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x) (\varphi - h). \quad (119)$$

Now, owing to the previous results $\varphi \in C([0, T], L^2(\Omega))$, the function $\bar{\varphi} : t \in [0, T] \mapsto \bar{\varphi}(t) \in [0, 1]$ is continuous. First, we show that there exists a constant $p > 0$, independent of t , that satisfies:

$$\bar{\varphi}(t) \geq p \quad \forall t \in [0, T]. \quad (120)$$

By contradiction, assume that there exists $t_1 \in (0, T]$ such that t_1 is the point where $\bar{\varphi}(t) = 0$, $t \in [0, T]$. From this, we deduce that $\varphi(t') = 0$ almost everywhere for $x \in \Omega$. We thus have that

$$\begin{aligned} \bar{\varphi}(t) &= \bar{\varphi}_0 + \int_0^t \int_{\Omega \setminus D} \chi_{\Omega \setminus D}(x) (\varphi - h) \\ &\geq \bar{\varphi}_0 + \int_0^t \int_{\Omega} (\varphi - h) \\ &\geq \bar{\varphi}_0 + \int_0^t \int_{\Omega} [(\varphi(m) - h) - (\varphi(t') - h)] dm \\ &\geq \bar{\varphi}_0 + \int_0^t \int_{\Omega} \varphi(m) dm \\ &\geq \bar{\varphi}_0 + \int_0^t \bar{\varphi}(m) dm. \end{aligned}$$

Applying Gronwall's Lemma, we get $\bar{\varphi}(t) \geq \bar{\varphi}_0 e^t > 0$ for all $t \in [0, T]$ (i.e., it is bounded below). However, we also have $\bar{\varphi}(t') = 0$, leading to a contradiction. Therefore, (120) is proved.

Additionally, let $d_0 = \frac{1}{2}p$. Then, we have:

$$|\Omega_1^t| \geq \frac{1}{2}p = e_0 > 0.$$

Now, suppose by contradiction that $|\Omega_1^t| < \frac{1}{2}p$. In

this case, we get:

$$\begin{aligned} p &\leq \bar{\varphi}(t) = \int_{\Omega_1^t} \varphi + \int_{\Omega \setminus \Omega_1^t} \varphi \\ &\leq \int_{\Omega_1^t} 1 + \int_{\Omega \setminus \Omega_1^t} \frac{p}{2} \\ &\leq |\Omega_1^t| + \frac{p}{2} \int_{\Omega \setminus \Omega_1^t} 1 \\ &\leq \frac{p}{2} + \frac{p}{2} |\Omega \setminus \Omega_1^t|. \end{aligned}$$

Thus, we arrive that $|\Omega \setminus \Omega_1^t| > 1$, which contradicts the assumption $|\Omega| = 1$. $\square$

Lemma 3.15 Consider that φ is a weak solution of the system (4)-(8) as defined in Definition 2.1, and assume that assumption (B) holds. Then, there exist constants $d_0 > 0$ and $e_0 > 0$, independent of t , such that:

$$|\Omega_1^t| \geq e_0 > 0, \quad (121)$$

to which $\Omega_1^t = \{x' \in \Omega; \varphi(x', t) \leq 1 - d_0\}$.

Proof: The verification of this lemma is identical to the verification of Lemma 4.1. $\square$

There are two main propositions that are essential for proving Theorem 5.1.

Proposition 3.16 Suppose the conditions of theorem 4.0.1 are satisfied. Then, for every $T_0 \in (0, T)$ there exists a constant $q > 0$, depending on both T_0 , and $\bar{\varphi}_0$, such that:

$$q \leq \varphi(x, t) \text{ for a.e. } x \in \Omega, t \in [T_0, T]. \quad (122)$$

Moreover, $T_0 = 0$ provided that there exists $\tilde{q} > 0$ such that:

$$\tilde{q} \leq \varphi_0 \text{ almost every where in } \Omega. \quad (123)$$

Proof: The verification of Proposition 5.1 involves proving that:

$$\|\ln(\varphi(\cdot, t))\|_{L^\infty(\Omega)} \leq R \quad \forall t \in [T_0, T].$$

Remark 3.17 Given that $\varphi \in L^\infty(0, T, H^2(\Omega)) \cap C([0, T], L^2(\Omega))$, it follows that $\varphi \in C([0, T], H^m(\Omega))$ for all $m < 2$, and $\varphi \in C([0, T], L^\infty(\Omega))$ if $l \leq 3$.

First, by Remark 5, consider $0 < \varphi(t)$ almost everywhere in Ω for all $t \in [0, T]$. Define

$$M_j(t) = \|\ln(\varphi(\cdot, t))\|_{L^\infty(\Omega)} \quad j = 1, 2, 3, \dots$$

Now, by choosing $j = 1$, and deriving a differential inequality of $M_j(t)$, we use the test function $\phi = \frac{1}{\varphi}$. This leads to the following:

$$\begin{aligned} \frac{d}{dt} M_1(t) &= \frac{d}{dt} \int_{\Omega} (-\ln(\varphi)) = - \left(\left(\varphi', \frac{1}{\varphi} \right) \right) \\ &= \int_{\Omega} \nabla \left(\frac{1}{\varphi} \right) (\nabla \varphi + \mu \nabla \theta) - \int_{\Omega} \frac{1}{\varphi} \chi_{\Omega \setminus D}(x) (\varphi - h). \end{aligned}$$

We compute

$$\frac{\chi_{\Omega \setminus D}(x) (\varphi - h)}{\varphi} \leq -\frac{\varphi - h}{\varphi} + \frac{-h}{\varphi} \leq -\frac{\varphi}{\varphi} = 1.$$

Then, we obtain

$$- \int_{\Omega} \frac{1}{\varphi} \chi_{\Omega \setminus D}(x) (\varphi - h) \leq E. \quad (124)$$

This can be verified by showing that:

$$|\nabla \ln(\varphi)|^2 = -\nabla \varphi \nabla \left(\frac{1}{\varphi} \right)$$

and

$$\frac{\nabla \ln(\varphi)}{\varphi} = -\nabla \left(\frac{1}{\varphi} \right).$$

Afterwards, by applying Young's inequality, we compute as follows:

$$\begin{aligned} \frac{d}{dt} M_1(t) &\leq - \int_{\Omega} |\nabla \ln(\varphi)|^2 - \int_{\Omega} \frac{\nabla \ln(\varphi)}{\varphi} \mu \nabla \theta + E \\ &\leq - \int_{\Omega} |\nabla \ln(\varphi)|^2 - \int_{\Omega} (1 - \varphi) \nabla \ln(\varphi) \nabla \theta + E \\ &\leq - \int_{\Omega} |\nabla \ln(\varphi)|^2 + \frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + R \int_{\Omega} |\nabla \theta|^2 + E \\ &\leq -\frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + R \int_{\Omega} |\nabla \theta|^2 + E \\ &\leq -\frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + R \int_{\Omega} |M * (1 - 2\varphi)|^2 + E \\ &\leq -\frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + R r_2^2 \int_{\Omega} |1 - 2\varphi|^2 + E \\ &\leq -\frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + E \int_{\Omega} |\varphi|^2 + R + E \\ &\leq -\frac{1}{2} \int_{\Omega} |\nabla \ln(\varphi)|^2 + E. \end{aligned} \quad (125)$$

Remark 3.18 Let $\Omega \subset \mathbb{R}^l$, where $l \in \mathbb{N}$ be a bounded domain. Suppose $\rho \in H^1(\Omega)$ and $\Omega' \subseteq \Omega$ with $|\Omega'| > 0$. Then, there exists a constant $E \geq 0$, dependent on Ω and Ω' , such that:

$$\left\| \rho - \frac{1}{|\Omega'|} \int_{\Omega'} \rho \right\|_{L^2(\Omega)} \leq E \frac{1}{|\Omega'|} \|\nabla \rho\|_{L^2(\Omega)}.$$

Consider d_0 , Ω_1^t and e_0 as defined in Lemma 4.1.

By remark 6, (2.11) and (4.3), we obtain:

$$\begin{aligned} \int_{\Omega} |\nabla \ln(\varphi)|^2 &\geq E|\Omega_1^t|^2 \int_{\Omega} \left| \ln(\varphi) - \frac{1}{|\Omega_1^t|} \int_{\Omega_1^t} \ln(\varphi) \right|^2 \\ &\geq E|\Omega_1^t|^2 \left[\int_{\Omega} |\ln(\varphi)|^2 - \frac{2}{|\Omega_1^t|} \int_{\Omega} \ln(\varphi) \int_{\Omega_1^t} \ln(\varphi) \right] \\ &\quad + E|\Omega_1^t|^2 \left[\int_{\Omega} \frac{1}{|\Omega_1^t|^2} \left(\int_{\Omega_1^t} \ln(\varphi) \right)^2 \right] \\ &\geq E|\Omega_1^t|^2 \left(\int_{\Omega} |\ln(\varphi)|^2 - \frac{2}{|\Omega_1^t|} \int_{\Omega} \ln(\varphi) \int_{\Omega_1^t} \ln(\varphi) \right) \\ &\geq Ee_0^2 \left(\left(\int_{\Omega} |\ln(\varphi)| \right)^2 + \frac{2}{e_0} \int_{\Omega} |\ln(\varphi)| \int_{\Omega_1^t} \ln(d_0) \right) \\ &\geq E(M_1(t))^2 - EM_1(t) \\ &\geq E(M_1(t))^2 - E. \end{aligned} \quad (126)$$

Now, combining (126) and (127), we deduce:

$$\frac{d}{dt} M_1(t) \leq -\frac{E}{2} (M_1(t))^2 + \frac{E}{2} + E \leq -E_1 (M_1(t))^2 + E_2.$$

According to [30], we can show that for all $T_0 \in [0, T]$, there exists A_1 dependent on $\bar{\varphi}_0$ and T_0 and independent of M_0 , such that:

$$M_1(t) \leq A_1 \quad \forall t \in [T_0, T]. \quad (127)$$

Likewise, we derive a differential inequality for M_j .

$$\begin{aligned} \frac{d}{dt} M_j(t) &= \frac{d}{dt} \left(\int_{\Omega} (-\ln(\varphi))^j \right) \\ &= \int_{\Omega} j \left(\left(\varphi', \frac{(-\ln(\varphi))^{j-1}}{\varphi} \right) \right) \left(-\frac{1}{j} \right) \left(\int_{\Omega} (-\ln(\varphi))^j \right)^{\frac{1}{j-1}} \\ &= M_j^{j-1} \int_{\Omega} \nabla \left(\frac{(-\ln(\varphi))^{j-1}}{\varphi} \right) (\nabla \varphi + \mu \nabla \theta) \\ &\quad - M_j^{j-1} \int_{\Omega} \chi_{\Omega \setminus D}(x) (\varphi - h) \frac{(-\ln(\varphi))^{j-1}}{\varphi}. \end{aligned}$$

Owing to Holder's inequality with conjugates j and $\frac{j}{j-1}$, we obtain the last term:

$$\begin{aligned} M_{j-1}^{j-1} &= \int_{\Omega} (-\ln(\varphi))^{j-1} \\ &\leq \left(\int_{\Omega} 1 \right)^{\frac{1}{j}} \left(\int_{\Omega} (-\ln(\varphi))^{(j-1)\frac{j}{j-1}} \right)^{\frac{j-1}{j}} \\ &\leq |\Omega|^{\frac{1}{j}} M_j^{j-1}. \end{aligned}$$

Then, using (4.10), we get:

$$\begin{aligned} &- M_j^{j-1} \int_{\Omega} \chi_{\Omega \setminus D}(x) (\varphi - h) \frac{(-\ln(\varphi))^{j-1}}{\varphi} \\ &\leq E M_j^{j-1} \int_{\Omega} (-\ln(\varphi))^{j-1} \\ &\leq E M_j^{j-1} M_{j-1}^{j-1} \\ &\leq E M_j^{j-1} |\Omega|^{\frac{1}{j}} M_j^{j-1} \leq E, \end{aligned} \quad (128)$$

where E is independent of j .

Next, by using (128) and [16], we can show the below differential inequality:

$$\frac{d}{dt} M_j \leq -E_3 \frac{1}{j^2} (M_j)^2 + E_4 A_1 M_1 + E_5 j,$$

where A_1 is as in (127), from which we deduce that $\forall \bar{T} \in (0, T]$, the following holds:

$$\sup_{t \geq \bar{T}} M_j(t) \leq R_1(\bar{T}) j^3 \quad \forall j \in [1, +\infty),$$

where $R_1(\bar{T})$ is independent of the initial values $M_j(0)$ and j , and is decreasing on $(0, +\infty)$. Additionally, $\bar{T} R_1(\bar{T})$ increases as $\bar{T}$ becomes large. Lastly, by [16], we conclude that for all $T_0 > 0$, there exists $R > 0$ such that:

$$M_j(t) \leq R \quad \forall t \geq T_0,$$

where R depends on T_0 and $\bar{\varphi}_0$. Taking the limit as $j \rightarrow \infty$, we deduce:

$$\|\ln(\varphi(\cdot, t))\|_{L^\infty(\Omega)} \leq R \quad \forall t \in [T_0, T].$$

□

Proposition 3.19 Assume that the conditions of Theorem 4.1 are satisfied. Then, for every $T_0 \in (0, T)$, there exists a constant $q > 0$, dependent on both T_0 and $\bar{\varphi}_0$, such that:

$$\varphi(x, t) \leq 1 - q \quad \text{a.e. } x \in \Omega \text{ and } t \in [T_0, T]. \quad (129)$$

Furthermore, if there exists $\tilde{q}$ such that

$$\varphi_0 \leq 1 - \tilde{q}, \quad (130)$$

then $T_0 = 0$.

Proof: The proof of Proposition 4.2 follows similarly to that of Proposition 4.1, differing only in the substitution $\Phi = 1 - \varphi$. □

As a result, by combining Propositions 4.1 and 4.2, we deduce the proof of Theorem 4.1.

4 Conclusion

The objectives of this paper were to study the existence, uniqueness, regularity, and separation properties of the nonlocal Cahn-Hilliard equation with a fidelity term and singular nonlinearities. Uniqueness of the solution was established by assuming the existence of two solutions, multiplying equations (??)–(??) by a test function $\varphi \in H^1(\Psi)$, and analyzing the difference with respect to the initial state. From these computations and the equality of the initial conditions, we concluded that the solution is unique.

Existence of a weak solution was proved by approximating the singular logarithmic nonlinearities with more regular functions. Using the Faedo-Galerkin method, we derived a priori estimates for the approximate solutions, established their existence, and passed to the limit to obtain a solution to the original problem. Regularity was demonstrated through four key lemmas, which were instrumental in proving that $\varphi \in L^\infty(0, T; H^2(\Omega))$ in the main theorem.

We also proved the separation property, showing that $\bar{\varphi}(t)$ remains bounded between 0 and 1, and established the existence of a global (in time) unique solution.

An interesting challenge in the nonlocal Cahn-Hilliard model is the development of numerical simulations to evaluate its efficiency compared to the classical model. In particular, the question arises: “What advantages does the nonlocal Cahn-Hilliard model offer over the classical one in practical applications?” In future work, we aim to address this by developing a convergent numerical scheme for the proposed model and performing simulations to highlight the potential benefits of the nonlocal approach.

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